

Matrix-analytic solution of second order Markov fluid models by using matrix-quadratic equations

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MAM 9, June. 28, 2016, Budapest, Hungary

Outline

There are efficient numerical methods for regular (first order) Markov fluid models.

Can we use them for second order Markov fluid models?

- 1 *Introduction*
- 2 *Model description*
- 3 *Model analysis*
- 4 *Main theorem*
- 5 *Boundary behavior*
- 6 *Final remarks*

Fluid models

Fluid models

- Fluid flows in to and out from an infinite buffer.
- Fluid flow is modulated by a background Markov chain.
 - first order: fluid level changes at a constant rate

$$\mathcal{Z}(t + \Delta) - \mathcal{Z}(t) = r_i \Delta \quad \text{in state } i.$$

- second order: normal distributed fluid increment

$$\mathcal{Z}(t + \Delta) - \mathcal{Z}(t) = \mathcal{N}(r_i \Delta, \sigma_i^2 \Delta) \quad \text{in state } i.$$

First order fluid models

Nice special cases allow symbolic solutions (Anick, Mitra, Sondhi).

Numerical methods

- *Spectral decomposition based methods (Kulkarni)*
- *Ricatti equation based solutions (Ahn-Ramaswami, Soares-Latouche)*
- *Quadratic matrix equation based solution (Ramaswami)*

Second order fluid models

Numerical methods

- Spectral decomposition based methods (Karandikar-Kulkarni)
- Transformation to first order differential equation with larger state space (Kulkarni)
- Quadratic matrix equation based solution

Second order Markov Fluid model

Main characterization of the stochastic processes:

- infinite buffer with lower boundary at level 0.
- $\mathcal{Z}(t)$: fluid level process.
- fluid increment process is characterized by
 - $\mathcal{X}(t)$: modulating CTMC with state space $\mathcal{S} = \{1, \dots, L\}$ and generator \mathbf{Q} ,
 - fluid rates, $r_j, j \in \{1, \dots, L\}$ described by diagonal matrix $\mathbf{R} = \text{diag}(r_1, \dots, r_L)$.
 - variance parameter, $\sigma_j, j \in \{1, \dots, L\}$ described by diagonal matrix $\mathbf{S} = \text{diag}(\sigma_1^2/2, \dots, \sigma_L^2/2)$

Second order Markov Fluid model

Performance measures of interest:

- stationary fluid level distribution

$$f_i(x) = \lim_{t \rightarrow \infty} \frac{d}{dx} P(\mathcal{X}(t) < x, \mathcal{Z}(t) = i),$$

- stationary buffer empty probability

$$p_i = \lim_{t \rightarrow \infty} P(\mathcal{X}(t) = 0, \mathcal{Z}(t) = i).$$

Analytical description

$f(x)$ satisfies the following differential and boundary equations

$$\frac{d}{dx}f(x)\mathbf{R} - \frac{d^2}{dx^2}f(x)\mathbf{S} = f(x)\mathbf{Q}, \quad (1)$$

$$f(0)\mathbf{R} - f'(0)\mathbf{S} = p\mathbf{Q}, \quad (2)$$

where $f'(0) = \left. \frac{d}{dx}f(x) \right|_{x=0}$.

State classification

The states in \mathcal{S} are divided into

- first order states with $\sigma_i^2 = 0$ and
- second order states ($\sigma_i^2 > 0$)
 - with *reflecting* boundary.
 - with *absorbing* boundary.

Properties

- In first order states:

$$p_i > 0, \forall i : r_i < 0 \text{ and } p_i = 0, \forall i : r_i > 0.$$

- In second order

- *reflecting* states:

$$p_i = 0, \forall i : \sigma_i^2 > 0,$$

- *absorbing* states:

$$f_i(0) = 0, \forall i : \sigma_i^2 > 0.$$

State classification

The state space \mathcal{S} is partitioned according to the sign of the rates and variances as follows:

- $\mathcal{S}^+ = \{i : r_i > 0, \sigma_i^2 = 0\}$, $\mathcal{S}^- = \{i : r_i < 0, \sigma_i^2 = 0\}$,
- $\mathcal{S}^{\sigma+} = \{i : r_i > 0, \sigma_i^2 > 0\}$, $\mathcal{S}^{\sigma-} = \{i : r_i < 0, \sigma_i^2 > 0\}$.

The set of states is decomposed as

$$\mathcal{S} = \mathcal{S}^+ \cup \mathcal{S}^{\sigma+} \cup \mathcal{S}^{\sigma-} \cup \mathcal{S}^- = \mathcal{S}^\bullet \cup \mathcal{S}^-, \text{ where}$$

$$\mathcal{S}^\bullet = \mathcal{S}^+ \cup \mathcal{S}^{\sigma+} \cup \mathcal{S}^{\sigma-}.$$

We assume that the states are numbered according to the $\mathcal{S}^+, \mathcal{S}^{\sigma+}, \mathcal{S}^{\sigma-}, \mathcal{S}^-$ order of subsets.

We exclude $r_i = 0$!!!

Solution of the differential equation

Similar to first order models, $f(x)$ can be expressed in a matrix-exponential form (Karandikar-Kulkarni)

$$f(x) = \pi e^{\mathbf{K}x} [I \quad \Psi], \quad (3)$$

where

- π is a row vector of size $|\mathcal{S}^\bullet|$,
- the size of \mathbf{K} is $|\mathcal{S}^\bullet| \times |\mathcal{S}^\bullet|$
- and the size of Ψ is $|\mathcal{S}^\bullet| \times |\mathcal{S}^-|$.

It remains to solve

- matrices \mathbf{K} and Ψ ,
- vector π ,
- and the vector of probability masses at level 0 p .

Quadratic equations

Substituting (3) into the differential equation (1) gives

$$\begin{aligned} KR_{\bullet} - K^2 S_{\bullet} &= Q_{\bullet\bullet} + \Psi Q_{-\bullet}, \\ K\Psi R_{-} - \underbrace{K^2 \Psi S_{-}}_{\mathbf{0}} &= Q_{\bullet-} + \Psi Q_{--}, \end{aligned}$$

where $S_{-} = \mathbf{0}$ has been exploited.

Our goal is to transform this set of quadratic equations into a single one of size $|\mathcal{S}|$ with proper signs of the coefficients.

Main theorem

Theorem

The minimal non-negative solution of the matrix-quadratic equation $\mathbf{F} + \mathbb{R}\mathbf{L} + \mathbb{R}^2\mathbf{B} = \mathbf{0}$ defined by the QBD with forward, local and backward matrix blocks

$$\mathbf{F} = \begin{bmatrix} \hat{\mathbf{Q}}_{..} + \hat{\mathbf{l}}_{.} + \hat{\mathbf{S}}_{.} & \hat{\mathbf{Q}}_{.-} \\ \mathbf{0} & \mathbf{0} \end{bmatrix}, \mathbf{L} = \begin{bmatrix} -\hat{\mathbf{l}}_{.} - 2\hat{\mathbf{S}}_{.} & \mathbf{0} \\ \hat{\mathbf{Q}}_{-.} & \hat{\mathbf{Q}}_{--} - \mathbf{I}_{-} \end{bmatrix},$$

$$\mathbf{B} = \begin{bmatrix} \hat{\mathbf{S}}_{.} & \mathbf{0} \\ \mathbf{0} & \mathbf{I}_{-} \end{bmatrix} \text{ is } \mathbb{R} = \begin{bmatrix} \hat{\mathbf{K}} + \mathbf{I}_{.} & \hat{\mathbf{\Psi}} \\ \mathbf{0} & \mathbf{0} \end{bmatrix}.$$

Main theorem

Where the notations with $\hat{\cdot}$ denotes properly scaled quantities

$$\mathbf{C}_\bullet = \begin{bmatrix} \mathbf{R}_+ & & \\ & \mathbf{R}_{\sigma+} & \\ & & -\mathbf{R}_{\sigma-} \end{bmatrix} \text{ and } \mathbf{C}_- = -\mathbf{R}_-,$$

$$c > \max \left(\max_{i \in S^+} \frac{-q_{ii}}{r_i}, \max_{i \in S^{\sigma-} \cup S^{\sigma+}} \frac{-r_i + \sqrt{r_i^2 - 2\sigma_i^2 q_{ii}}}{\sigma_i^2} \right),$$

$$\hat{\mathbf{K}} = \frac{1}{c} \mathbf{C}_\bullet^{-1} \mathbf{K} \mathbf{C}_\bullet, \quad \hat{\Psi} = \frac{1}{c} \mathbf{C}_\bullet^{-1} \Psi \mathbf{C}_-, \quad \hat{\mathbf{S}}_\bullet = c \mathbf{C}_\bullet^{-1} \mathbf{S}_\bullet, \quad \hat{\mathbf{Q}} = \frac{1}{c} \mathbf{C}^{-1} \mathbf{Q},$$

$$\text{and } \hat{\mathbf{I}}_\bullet = \mathbf{C}_\bullet^{-1} \mathbf{R}_\bullet = \begin{bmatrix} \mathbf{I}_+ & & \\ & \mathbf{I}_{\sigma+} & \\ & & -\mathbf{I}_{\sigma-} \end{bmatrix}.$$

Main theorem

The elements of the proof:

- Simple substitution provides the identity with the differential equation.
- The scaling ensures that \mathbf{B} , \mathbf{L} and \mathbf{F} are proper QBD matrix blocks.
- The QBD solution ensured the (minimal non-negative) solution with the proper eigenvalues.

Reflecting boundary at second order states

In all second order states then $p_{\bullet} = 0$

Inserting the matrix-exponential solution into (2) leads to

$$\pi \mathbf{R}_{\bullet} - \pi \mathbf{K} \mathbf{S}_{\bullet} = p_{-} \mathbf{Q}_{-\bullet}, \quad (4)$$

$$\pi \Psi \mathbf{R}_{-} = p_{-} \mathbf{Q}_{--}, \quad (5)$$

since $\mathbf{S}_{-} = 0$.

From (5) and (4)

$$\begin{aligned} \pi(\mathbf{R}_{\bullet} - \mathbf{K} \mathbf{S}_{\bullet} - \Psi \mathbf{R}_{-} (\mathbf{Q}_{--})^{-1} \mathbf{Q}_{-\bullet}) &= 0, \\ p_{-} &= \pi \Psi \mathbf{R}_{-} (\mathbf{Q}_{--})^{-1}, \end{aligned}$$

where $\mathbf{R}_{-} (\mathbf{Q}_{--})^{-1}$ is a non-negative matrix.

The normalization condition ($\int_x f(x) \mathbb{1} + p_{-} \mathbb{1} = 1$), is

$$\pi \left((-\mathbf{K})^{-1} \begin{bmatrix} \mathbf{I} & \Psi \end{bmatrix} \mathbb{1} + \mathbf{R}_{-} (\mathbf{Q}_{--})^{-1} \mathbb{1} \right) = 1.$$

Absorbing boundary at second order states

$f_{\sigma_+}(0) = 0$ and $f_{\sigma_-}(0) = 0$.

$f(0) = \pi [I \quad \Psi]$ implies that $\pi_{\sigma_+} = 0$ and $\pi_{\sigma_-} = 0$.

Substituting it into (2) gives

$$f(0)\mathbf{R} = [\pi_+ \mathbf{R}_+ \quad 0 \quad 0 \quad \pi_+ \Psi_{+-} \mathbf{R}_-],$$

$$f'(0)\mathbf{S} = [0 \quad \pi_+ \mathbf{K}_{+,\sigma_+} \mathbf{S}_{\sigma_+} \quad \pi_+ \mathbf{K}_{+,\sigma_-} \mathbf{S}_{\sigma_-} \quad 0],$$

since $\mathbf{S}_+ = \mathbf{S}_- = 0$.

Absorbing boundary at second order states

For the partitioned vectors and block matrices (2) can be rewritten as

$$\begin{aligned}\pi_+ \mathbf{R}_+ &= \rho_{\sigma_+} \mathbf{Q}_{\sigma_+,+} + \rho_{\sigma_-} \mathbf{Q}_{\sigma_-,+} + \rho_- \mathbf{Q}_{-,+}, \\ -\pi_+ \mathbf{K}_{+,\sigma_+} \mathbf{S}_{\sigma_+} &= \rho_{\sigma_+} \mathbf{Q}_{\sigma_+,\sigma_+} + \rho_{\sigma_-} \mathbf{Q}_{\sigma_-,\sigma_+} + \rho_- \mathbf{Q}_{-,\sigma_+}, \\ -\pi_+ \mathbf{K}_{+,\sigma_-} \mathbf{S}_{\sigma_-} &= \rho_{\sigma_+} \mathbf{Q}_{\sigma_+,\sigma_-} + \rho_{\sigma_-} \mathbf{Q}_{\sigma_-,\sigma_-} + \rho_- \mathbf{Q}_{-,\sigma_-}, \\ \pi_+ \boldsymbol{\Psi}_{+-} \mathbf{R}_- &= \rho_{\sigma_+} \mathbf{Q}_{\sigma_+,-} + \rho_{\sigma_-} \mathbf{Q}_{\sigma_-,-} + \rho_- \mathbf{Q}_{-,-},\end{aligned}$$

which gives the linear system

$$\begin{bmatrix} \pi_+ & \rho_{\sigma_+} & \rho_{\sigma_-} & \rho_- \\ -\mathbf{R}_+ & \mathbf{K}_{+,\sigma_+} \mathbf{S}_{\sigma_+} & \mathbf{K}_{+,\sigma_+} \mathbf{S}_{\sigma_-} & -\boldsymbol{\Psi}_{+-} \mathbf{R}_- \\ \mathbf{Q}_{\sigma_+,+} & \mathbf{Q}_{\sigma_+,\sigma_+} & \mathbf{Q}_{\sigma_+,\sigma_-} & \mathbf{Q}_{\sigma_+,-} \\ \mathbf{Q}_{\sigma_-,+} & \mathbf{Q}_{\sigma_-,\sigma_+} & \mathbf{Q}_{\sigma_-,\sigma_-} & \mathbf{Q}_{\sigma_-,-} \\ \mathbf{Q}_{-,+} & \mathbf{Q}_{-,\sigma_+} & \mathbf{Q}_{-,\sigma_-} & \mathbf{Q}_{-,-} \end{bmatrix} = 0.$$

Absorbing boundary at second order states

Whose normalization condition is

$$[\pi_+ \quad \rho_{\sigma+} \quad \rho_{\sigma-} \quad \rho_-] \cdot \begin{bmatrix} (-\mathbf{K})_+^{-1} [I \quad \Psi] \mathbb{1} \\ \mathbb{1} \\ \mathbb{1} \\ \mathbb{1} \end{bmatrix} = \mathbf{1}.$$

Final remarks

Summary

- 1 It was possible to recycle the methodology developed for first order fluid models for the analysis of second order fluid models.
- 2 The proof is based on known properties of QBD quadratic matrix equations.
- 3 Similar to first order fluid models, based on \mathbf{K} and Ψ the boundary conditions are obtained from the solution of a linear system.

Plans

- 1 analysis of further second order fluid models.